BOUNDARY VALUE PROBLEMS (I)

• Solving a two-point boundary value problem

$$\frac{d^2 y}{dx^2} = g \qquad \text{for } x \in (0, 2\pi)$$
$$y(0) = y(2\pi) = 0$$

- Finite-difference approximation:
 - Second–order Central Difference formula for the interior nodes:

$$\frac{y_{j+1} - 2y_j + y_{j-1}}{h^2} = g_j \text{ for } j = 1, \dots, N$$

where $h = \frac{2\pi}{N+1}$ and $x_j = jh$
- Endpoint nodes:
 $y_0 = 0 \implies y_2 - 2y_1 = h^2 g_1$
 $y_{N+1} = 0 \implies -2y_N + y_{N-1} = h^2 g_N$

- Tridiagonal algebraic system

Numerical Methods for ODEs

BOUNDARY VALUE PROBLEMS (II)

PART II

Review of numerical methods for

Ordinary Differential Equations

• Solving a two-point boundary value problem

$$\frac{d^2 y}{dx^2} = g \qquad \text{for } x \in (0, 2\pi)$$
$$\frac{dy}{dx}(0) = \frac{dy}{dx}(2\pi) = 0$$

- Finite-difference approximation:
 - Second–order Central Difference formula for the interior nodes:

$$\frac{y_{j+1} - 2y_j + y_{j-1}}{h^2} = g_j \text{ for } j = 1, \dots, N$$

- First-order Forward/Backward Difference formulae to re-express endpoint values: $y_1 - y_0 = 0 \implies y_0 = y_1$

$$\frac{1}{h} = 0 \implies y_0 = y_1$$
$$\frac{y_{N+1} - y_N}{h} = 0 \implies y_{N+1} = y_N$$

First-order only --- degraded accuracy!

- Tridiagonal algebraic system — Where is the problem?

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BOUNDARY VALUE PROBLEMS (III)

• In order to retain second–order accuracy in the approximation of the Neumann problem need to use higher-order formulae at endpoints, e.g.

$$y'_0 = \frac{-y_2 + 4y_1 - 3y_0}{2h} = 0 \implies y_0 = \frac{1}{3}(-y_2 + 4y_1)$$

• The first row thus becomes

$$\frac{2}{3}y_2 - \frac{2}{3}y_1 = h^2 g_1$$

Second-order accuracy recovered!

BOUNDARY VALUE PROBLEMS (IV)

- Compact Stencils stencils based on three grid points only: $\{x_{j+1}, x_j, x_{j-1}\}$ at the j th node
- Is is possible to obtain higher (then second) order of accuracy on compact stencils? YES!
- Consider the central difference approximation to the equation $\frac{d^2y}{dy^2} = g$

$$\frac{y_{j+1} - 2y_j + y_{j-1}}{h^2} - \frac{h^2}{12}y_j^{(iv)} + O(h^4) = g_j$$

• Re-express the error term $\frac{h^2}{12}y_i^{(iv)}$ using the equation in question:

$$\frac{h^2}{12}y_j^{(iv)} = \frac{h^2}{12}g_j'' = \frac{h^2}{12} \left[\frac{g_{j+1} - 2g_j + g_{j-1}}{h^2} - \frac{h^2}{12}g_j^{(iv)} + O(h^4)\right]$$

• Inserting into the original finite-difference equation:

$$\frac{y_{j+1} - 2y_j + y_{j-1}}{h^2} = g_j + \frac{g_{j+1} - 2g_j + g_{j-1}}{12} + O(h^4)$$

• Slight modification of the RHS \implies fourth—order accuracy!!!

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INITIAL VALUE PROBLEMS — GENERAL REMARKS

• Consider the following Cauchy problem :

$$\frac{dy}{dt} = f(y,t) \text{ with } y(t_0) = y_0$$

The independent variable t is usually referred to as time .

- Equations with higher–order derivatives can be reduced to systems of first–order equations
- Generalizations to systems of ODEs straightforward
- When the RHS function doesn't depend on y, i.e., f(y,t) = f(t), solution obtained via quadrature
- Assume uniform time-steps (*h* is constant)

BOUNDARY VALUE PROBLEMS (V)

- Compact Finite Difference Schemes drawbacks:
 - need to be tailored to the specific equation solved
 - can get fairly complicated for more complex equations

Numerical Methods for ODEs

INITIAL VALUE PROBLEMS — CHARACTERIZATION OF INTEGRATION METHODS

• ACCURACY — unlike in the Boundary Value Problems, there is no terminal condition and approximation errors may accumulate in time; consequently, a relevant characterization of accuracy is provided by the global error

 $(\text{global error}) = (\text{local error}) \times (\# \text{ of time steps}),$

rather than the local error.

• **STABILITY** — unlike in the Boundary Value Problems, where boundedness of the solution at final time is enforced via a suitable terminal condition, in Initial Value Problems there is a priori no guarantee that the solution will remain bounded.

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INITIAL VALUE PROBLEMS — MODEL PROBLEM

• Stability of various numerical schemes is usually analyzed by applying these schemes to the following linear model:

$$\frac{dy}{dt} = \lambda y = (\lambda_r + i\lambda_i)y$$
 with $y(t_0) = y_0$

which is stable when $\lambda_r <= 0$.

• Exact solution:
$$y(t) = y_0 e^{\lambda t} = \left(1 + \lambda h + \frac{\lambda^2 h^2}{2} + \frac{\lambda^3 h^3}{6} + \dots\right) y_0$$

• Motivation — consider the following advection-diffusion PDE:

$$\frac{\partial u}{\partial t} + c\frac{\partial u}{\partial x} - a\frac{\partial^2 u}{\partial x^2} = 0$$

Taking Fourier transform yields (*k* is the wavenumber):

$$\frac{d\hat{u}_k}{dt} + c\,ik\,\hat{u}_k + ak^2\,\hat{u}_k = 0$$

where

- the real term $ak^2 \hat{u}_k$ represents diffusion
- the imaginary term $cik\hat{u}_k$ represents advection

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INITIAL VALUE PROBLEMS — EXPLICIT EULER SCHEME (II)

• Local error analysis:

$$y_{n+1} = (1 + \lambda h)y_n + [O(h^2)]$$

• Global error analysis:

(global error) =
$$Ch^2 \cdot N = Ch^2 \cdot \frac{T}{h} = C'h$$

Thus, the scheme is

- locally second-order accurate
- globally (over the interval $[t_0, t_0 + Nh]$) first-order accurate

INITIAL VALUE PROBLEMS — EXPLICIT EULER SCHEME (I)

• Consider a Taylor series expansion

$$y(t_{n+1}) = y(t_n) + hy'(t_n) + \frac{h^2}{2}y''(t_n) + \dots$$

Using the ODE we obtain

$$y' = \frac{dy}{dt} = f$$
$$y'' = \frac{dy'}{dt} = \frac{df}{dt} = f_t + f_f$$

• Neglecting terms proportional to second and higher powers of *h* yields the Explicit Euler Method

 $y_{n+1} = y_n + hf(y_n, t_n)$

• Retaining higher–order terms is inconvenient, as it requires differentiation of *f* and does not lead to schemes with desirable stability properties.

Numerical Methods for ODEs

Numerical Methods for ODEs

INITIAL VALUE PROBLEMS — EXPLICIT EULER SCHEME (III)

• Stability (for the model problem)

$$y_{n+1} = y_n + \lambda h y_n = (1 + \lambda h) y_n$$

Thus, the solution after n time steps

$$y_n = (1 + \lambda h)^n y_0 \triangleq \sigma^n y_0 \implies \sigma = 1 + \lambda h$$

For large *n*, the numerical solution remains stable iff

$$\sigma \leq 1 \implies (1+\lambda_r h)^2 + (\lambda_i h)^2 \leq 1$$

- conditionally stable for real $\boldsymbol{\lambda}$
- unstable stable for imaginary $\boldsymbol{\lambda}$

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INITIAL VALUE PROBLEMS — Implicit Euler Scheme (I)

- Implicit Schemes based on approximation of the RHS that involve $f(y_{n+1},t)$, where y_{n+1} is the unknown to be determined
- Implicit Euler Scheme obtained by neglecting second and higher–order terms in the expansion:

$$y(t_n) = y(t_{n+1}) - hy'(t_{n+1}) + \frac{h^2}{2}y''(t_{n+1}) - \dots$$

Upon substitution $\frac{dy}{dt}\Big|_{t_{n+1}} = f(y_{n+1}, t_{n+1})$ we obtain

 $y_{n+1} = y_n + hf(y_{n+1}, t_{n+1})$

The scheme is

- locally second-order accurate
- globally (over the interval $[t_0, t_0 + Nh]$) first-order accurate

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INITIAL VALUE PROBLEMS — CRANK–NICHOLSON SCHEME (I)

• Obtained by approximating the formal solution of the ODE $y_{n+1} = y_n + \int_{t_n}^{t_{n+1}} f(y,t) dt$ using the trapezoidal quadrature:

$$y_{n+1} = y_n + \frac{h}{2} [f(y_n, t_n) + f(y_{n+1}, t_{n+1})]$$

The scheme is

- locally third-order accurate
- globally (over the interval $[t_0, t_0 + Nh]$) second-order accurate
- Stability (for the model problem):

$$y_{n+1} = y_n + \frac{\lambda h}{2} (y_{n+1} + y_n) \implies y_{n+1} = \left(\frac{1 + \frac{\lambda h}{2}}{1 - \frac{\lambda h}{2}}\right) y_n$$
$$y_{n+1} = \left(\frac{1 + \frac{\lambda h}{2}}{1 - \frac{\lambda h}{2}}\right)^n y_0 \triangleq \sigma^n y_0 \implies \sigma = \frac{1 + \frac{\lambda h}{2}}{1 - \frac{\lambda h}{2}}$$
$$|\sigma| \le 1 \implies \Re(\lambda h) \le 0$$

Stable for all model ODEs with stable solutions

INITIAL VALUE PROBLEMS — IMPLICIT EULER SCHEME (II)

• Stability (for the model problem):

$$y_{n+1} = y_n + \lambda h y_{n+1} \implies y_{n+1} = (1 - \lambda h)^{-1} y_n$$
$$y_{n+1} = \left(\frac{1}{1 - \lambda h}\right)^n y_0 \triangleq \sigma^n y_0 \implies \sigma = \frac{1}{1 - \lambda h}$$
$$|\sigma| \le 1 \implies (1 - \lambda r h)^2 + (\lambda_i h)^2 \ge 1$$

Implicit Euler scheme is thus stable for

- all stable model problems
- most unstable model problems
- When solving systems of ODEs of the form $\mathbf{y} = \mathcal{A}(t)\mathbf{y}$, each implicit step requires solution of an algebraic system: $\mathbf{y}_{n+1} = (I h\mathcal{A})^{-1}\mathbf{y}_n$
- Implicit schemes are generally hard to implement for nonlinear problems

Numerical Methods for ODEs

INITIAL VALUE PROBLEMS — LEAPFROG SCHEME (I)

• Leapfrog as an example of a two-step method :

 $y_{n+1} = y_{n-1} + 2h\lambda y_n$

• Characteristic equation for the amplification factor $(y_n = \sigma^n y_0)$

$$\sigma^2 - 2h\lambda\sigma - 1 = 0$$

where roots give the amplification factors:

$$\sigma_1 = \lambda h + \sqrt{1 + \lambda^2 h^2} \simeq 1 + \lambda h + \frac{\lambda^2 h^2}{2} + \dots = e^{\lambda h} + O(h^3)$$

$$\sigma_2 = \lambda h - \sqrt{1 + \lambda^2 h^2} \simeq -(1 - \lambda h + \frac{\lambda^2 h^2}{2} - \dots) = -e^{-\lambda h} + O(h^3)$$

Thus, the scheme is

- locally third-order accurate
- globally (over the interval $[t_0, t_0 + Nh]$) second-order accurate

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INITIAL VALUE PROBLEMS — LEAPFROG SCHEME (II)

• Stability for diffusion problems ($\lambda = \lambda_r$):

$$\sigma_1 = \lambda h + \sqrt{1 + \lambda_r^2 h^2} > 1$$
 for all $h > 0$

Thus the scheme is unconditionally unstable for diffusion problem!

• Stability for advection problems ($\lambda = i\lambda_i$):

$$\sigma_{1/2}^2 = 1$$
 (!!!) for $h < \frac{1}{|\lambda_i|}$

Thus the scheme is conditionally unstable and non-diffusive for advection problems!

• Question — analyze dispersive (i.e., related to $\arg(\sigma)$) errors of the leapfrog scheme.

Numerical Methods for ODEs

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INITIAL VALUE PROBLEMS — RUNGE-KUTTA METHODS (I)

• General form of a fractional step method :

 $y_{n+1} = y_n + \gamma_1 h k_1 + \gamma_2 h k_2 + \gamma_3 h k_3 + \dots$

where

- $k_{1} = f(y_{n}, t_{n})$ $k_{2} = f(y_{n} + \beta_{1}hk_{1}, t_{n} + \alpha_{1}h)$ $k_{3} = f(y_{n} + \beta_{2}hk_{1} + \beta_{3}hk_{2}, t_{n} + \alpha_{2}h)$
- Choose γ_i , β_i and α_i to match as many expansion coefficients as possible in

$$y(t_{n+1}) = y(t_n) + hy'(t_n) + \frac{h^2}{2}y''(t_n) + \frac{h^3}{6}y'''(t_n) \dots$$

$$y' = f$$

$$y'' = f_t + ff_y$$

$$y''' = f_{tt} + f_t f_y 2ff_{yt} + f^2 f_{yt} + f^2 f_{yy}$$

• Runge—Kutta methods are self-starting with fairly good stability and accuracy properties.

INITIAL VALUE PROBLEMS — MULTISTEP PROCEDURES

• General form of a multistep procedure :

$$\sum_{j=1}^{p} \alpha_{j} y_{n+j} = h \sum_{j=1}^{q} \beta_{j} f(y_{n+j}, t_{n+j})$$

with characteristic polynomials

$$\xi_p(z) = \alpha_p z^p + \alpha_{p-1} z^{p-1} + \dots + \alpha_0$$

$$\zeta_q(z) = \beta_q z^q + \beta_{q-1} z^{q-1} + \dots + \beta_0$$

- if p > q explicit scheme
- if $p \le q$ implicit scheme
- A (ξ, ζ) -procedure converges uniformly in [a, b], i.e., $\lim_{h\to 0} \max_{t_n \in [a,b]} |y_n - y(t_n)| = 0$ if:
 - the following consistency conditions are verified: $\xi(1) = 0$ and $\xi'(1) = \zeta(1)$ (*consistency condition*)
 - all roots of the polynomial $\xi(z)$ are such that $|z_i| \le 1$ and the roots with $|z_k| = 1$ are simple (*stability condition*)

Numerical Methods for ODEs

INITIAL VALUE PROBLEMS — RUNGE-KUTTA METHODS (II)

• **RK4** — an ODE workhorse:

$$y_{n+1} = y_n + \frac{h}{6}k_1 + \frac{h}{3}(k_2 + k_3) + \frac{h}{6}k_4$$

$$k_1 = f(y_n, t_n) \qquad \qquad k_2 = f(y_n + \frac{h}{2}k_1, t_{n+1/2})$$

$$k_3 = f(y_n + \frac{h}{2}k_2, t_{n+1/2}) \qquad \qquad k_4 = f(y_n + hk_3, t_{n+1})$$

• The amplification factor:

$$\sigma = 1 + \lambda h + \frac{\lambda^2 h^2}{2} + \frac{\lambda^3 h^3}{6} + \frac{\lambda^4 h^4}{24}$$

Thus, stability iff $\sigma \leq 1$

• Accuracy:

$$e^{\lambda h} = \sigma + O(h^5)$$

Thus, the scheme is

- locally fifth-order accurate
- globally (over the interval $[t_0, t_0 + Nh]$) fourth-order accurate

INITIAL VALUE PROBLEMS — RUNGE'S PRINCIPLE

• Let (k+1) be the local truncation error; denote Y(t,h) an approximation of the exact solution y(t) computed with the step size h; then at $t = t_0 + 2nh$:

$$y(t) - Y(t,h) \simeq C 2nh^{k+1} = C(t-t_0)h^k$$

 $y(t) - Y(t,2h) \simeq Cn(2h)^{k+1} = C(t-t_0)2^kh^k$

Subtracting:

$$Y(t,2h) - Y(t,h) \simeq C(t-t_0)(1-2^k)h^k$$

Thus we can obtain an estimate of the absolute error based on solution with two step–sizes only:

$$y(t) - Y(t,h) \simeq \frac{Y(t,h) - Y(t,2h)}{2^k - 1}$$

• Runge's principle is very useful for adaptive step size refinement