

WELCOME TO MATH 745 – – TOPICS IN NUMERICAL ANALYSIS

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REVIEW OF NUMERICAL DIFFERENTIATION — FINITE DIFFERENCE FORMULAE I

- Approximation of derivatives $\frac{df}{dx}$ on a discrete set of points x_0, x_1, \dots, x_N
- Definitions & Assumptions:
 - $f_j = f(x_j)$
 - uniform mesh with constant grid spacing $h = x_{j+1} - x_j$
(extensions to nonuniform grids are straightforward)
- Derivation of finite difference formulae is based on Taylor-series expansions of the following form:

$$\begin{aligned} f_{j+1} &= f_j + (x_{j+1} - x_j)f'_j + \frac{(x_{j+1} - x_j)^2}{2!}f''_j + \frac{(x_{j+1} - x_j)^3}{3!}f'''_j + \dots \\ &= f_j + hf'_j + \frac{h^2}{2}f''_j + \frac{h^3}{6}f'''_j + \dots \end{aligned}$$

FINITE DIFFERENCE FORMULAE (II) — FORWARD-DIFFERENCE FORMULA

- Rearrange the Taylor series expansion

$$\begin{aligned}f'_j &= \frac{f_{j+1} - f_j}{h} - \frac{h}{2}f''_j + \dots \\ &= \frac{f_{j+1} - f_j}{h} + O(h),\end{aligned}$$

$O(h^\alpha)$ denotes the contribution from all terms with powers of h greater or equal α .

- Neglecting $O(h)$, we obtain the **first order forward-difference formula** :

$$\left(\frac{\delta f}{\delta x}\right)_j = \frac{f_{j+1} - f_j}{h}$$

- Neglected term with the lowest power of h is the **leading-order error**
- Exponent of h in the leading-order error represents the **order of accuracy of the method**
- Here: $Err = -\frac{h}{2}f''_j$, hence this method is **first order accurate**

FINITE DIFFERENCE FORMULAE (III) — BACKWARD DIFFERENCE FORMULA

- Backward difference formula is obtained by expanding f_{j-1} about x_j and proceeding as before:

$$f'_j = \frac{f_j - f_{j-1}}{h} - \frac{h}{2} f''_j + \dots \implies \left(\frac{\delta f}{\delta x} \right)_j = \frac{f_j - f_{j-1}}{h}$$

FINITE DIFFERENCE FORMULAE (IV) — HIGHER ORDER FORMULAE

- Consider two expansions:

$$f_{j+1} = f_j + hf'_j + \frac{h^2}{2} f''_j + \frac{h^3}{6} f'''_j + \dots$$

$$f_{j-1} = f_j - hf'_j + \frac{h^2}{2} f''_j - \frac{h^3}{6} f'''_j + \dots$$

- Subtracting the second from the first:

$$f_{j+1} - f_{j-1} = 2hf'_j + \frac{h^3}{3} f'''_j + \dots$$

- Central Difference Formula**

$$f'_j = \frac{f_{j+1} - f_{j-1}}{h} - \frac{h^2}{6} f'''_j + \dots \implies \left(\frac{\delta f}{\delta x} \right)_j = \frac{f_{j+1} - f_{j-1}}{2h}$$

- Leading-order error is $\frac{h^2}{6} f'''_j$, thus the method is **second-order accurate**
- Manipulating four different Taylor series expansions one can obtain a **fourth-order central difference formula** :

$$\left(\frac{\delta f}{\delta x} \right)_j = \frac{-f_{j+2} + 8f_{j+1} - 8f_{j-1} + f_{j-2}}{12h}$$

FINITE DIFFERENCE FORMULAE (V) — APPROXIMATION OF THE SECOND DERIVATIVE

- Consider two expansions:

$$f_{j+1} = f_j + hf'_j + \frac{h^2}{2} f''_j + \frac{h^3}{6} f'''_j + \dots$$

$$f_{j-1} = f_j - hf'_j + \frac{h^2}{2} f''_j - \frac{h^3}{6} f'''_j + \dots$$

- Adding the two expansions

$$f_{j+1} + f_{j-1} = 2f_j + h^2 f''_j + \frac{h^4}{12} f^{iv}_j + \dots$$

- Central difference formula for the second derivative:

$$f''_j = \frac{f_{j+1} - 2f_j + f_{j-1}}{h^2} - \frac{h^2}{12} f^{iv}_j + \dots \implies \left(\frac{\delta^2 f}{\delta x^2} \right)_j = \frac{f_{j+1} - 2f_j + f_{j-1}}{h^2}$$

- Leading-order error is $\frac{h^2}{12} f^{iv}_j$, thus the method is **second-order accurate**

FINITE DIFFERENCE FORMULAE (VI) — TAYLOR TABLE

- A general method for choosing the coefficients of a finite difference formula to ensure the highest possible order of accuracy
- Example: consider a one-sided finite difference formula $\sum_{p=0}^2 \alpha_p f_{j+p}$, where the coefficients α_p , $p = 0, 1, 2$ are to be determined.
- Form an expression for the approximation error

$$f'_j - \sum_{p=0}^2 \alpha_p f_{j+p} = \varepsilon$$

and expand it about x_j in the powers of h

FINITE DIFFERENCE FORMULAE (VII) — TAYLOR TABLE

- Expansions can be collected in a **Taylor table**

	f_j	f'_j	f''_j	f'''_j
f'_j	0	1	0	0
$-a_0 f_j$	$-a_0$	0	0	0
$-a_1 f_{j+1}$	$-a_1$	$-a_1 h$	$-a_1 \frac{h^2}{2}$	$-a_1 \frac{h^3}{6}$
$-a_2 f_{j+2}$	$-a_2$	$-a_2(2h)$	$-a_2 \frac{(2h)^2}{2}$	$-a_2 \frac{(2h)^3}{6}$

- the leftmost column contains the terms present in the expression for the approximation error
 - the corresponding rows (multiplied by the top row) represent the terms obtained from expansions about x_j
 - columns represent terms with the same order in h — sums of columns are the contributions to the approximation error with the given order in h
- The coefficients α_p , $p = 0, 1, 2$ can now be chosen to cancel the contributions to the approximation error with the **lowest powers of h**

FINITE DIFFERENCE FORMULAE (VIII) — TAYLOR TABLE

- Setting the coefficients of the first three terms to zero:

$$\begin{cases} -a_0 - a_1 - a_2 = 0 \\ 1 - a_1 h - a_2 (2h) = 0 \\ -a_1 \frac{h^2}{2} - a_2 \frac{(2h)^2}{2} = 0 \end{cases} \implies a_0 = -\frac{3}{2h}, \quad a_1 = \frac{2}{h}, \quad a_2 = -\frac{1}{2h}$$

- The resulting formula:

$$\left(\frac{\delta f}{\delta x} \right)_j = \frac{-f_{j+2} + 4f_{j+1} - 3f_j}{2h}$$

- The approximation error — determined the evaluating the first column with non-zero coefficient:

$$\left(-a_1 \frac{h^3}{6} - a_2 \frac{(2h)^3}{6} \right) f_j''' = \frac{h^2}{3} f_j'''$$

The formula is thus **second order accurate**

FINITE DIFFERENCE FORMULAE (IX) — SUBTRACTIVE CANCELLATION ERRORS

- **Subtractive cancellation errors** — when comparing two numbers which are almost the same using **finite-precision arithmetic**, the relative round-off error is proportional to the inverse of the difference between the two numbers
- Thus, if the difference between the two numbers is decreased by an order of magnitude, the relative accuracy with which this difference may be calculated using **finite-precision arithmetic** is also decreased by an order of magnitude.
- Problems with finite difference formulae when $h \rightarrow 0$ — less of precision due to finite-precision arithmetic (**subtractive cancellation**), e.g., for double precision:

$$1.0000000000000001 - 1 \sim 0$$

FINITE DIFFERENCE FORMULAE (X) — COMPLEX STEP DERIVATIVE^a

- Consider the complex extension $f(z)$, where $z = x + iy$, of $f(x)$ and compute the complex Taylor series expansion

$$f(x_j + ih) = f_j + ihf'_j - \frac{h^2}{2}f''_j - i\frac{h^3}{6}f'''_j + O(h^4)$$

- Take **imaginary** part and divide by h

$$f'_j = \frac{\Im(f(x_j + ih))}{h} + \frac{h^2}{6}f'''_j + O(h^3) \implies \left(\frac{\delta f}{\delta x}\right)_j = \frac{\Im(f(x_j + ih))}{h}$$

- Note that the scheme is **second order accurate** — where is conservation of complexity?
- The method doesn't suffer from cancellation errors, is easy to implement and quite useful

^aJ. N. Lyness and C. B. Moler, “Numerical differentiation of analytical functions”, *SIAM J. Numer Anal* **4**, 202-210, (1967)

FINITE DIFFERENCE FORMULAE (XI) — PADÉ APPROXIMATION

- GENERAL IDEA — include in the finite-difference formula not only the **function values** , but also the values of the **function derivative** at the adjacent nodes, e.g.:

$$b_{-1}f'_{j-1} + f'_j + b_1f'_{j+1} - \sum_{p=-1}^1 \alpha_p f_{j+p} = \varepsilon$$

- Construct the **Taylor table** using the following expansions:

$$f_{j+1} = f_j + hf'_j + \frac{h^2}{2}f''_j + \frac{h^3}{6}f'''_j + \frac{h^4}{24}f_j^{(iv)} + \frac{h^5}{120}f_j^{(v)} + \dots$$

$$f'_{j+1} = f'_j + hf''_j + \frac{h^2}{2}f'''_j + \frac{h^3}{6}f_j^{(iv)} + \frac{h^4}{24}f_j^{(v)} + \dots$$

NOTE — need an expansion for the derivative and a higher order expansion for the function (more coefficient to determine)

FINITE DIFFERENCE FORMULAE (XII) — PADÉ APPROXIMATION

- The Taylor table

	f_j	f'_j	f''_j	f'''_j	$f_j^{(iv)}$	$f_j^{(v)}$
$b_{-1}f'_{j-1}$	0	b_{-1}	$b_{-1}(-h)$	$b_{-1}\frac{(-h)^2}{2}$	$b_{-1}\frac{(-h)^3}{6}$	$b_{-1}\frac{(-h)^4}{24}$
f'_j	0	1	0	0	0	0
$b_1f'_{j+1}$	0	b_1	b_1h	$b_1\frac{h^2}{2}$	$b_1\frac{h^3}{6}$	$b_1\frac{h^4}{24}$
$-a_{-1}f_{j-1}$	$-a_{-1}$	$-a_{-1}(-h)$	$-a_{-1}\frac{(-h)^2}{2}$	$-a_{-1}\frac{(-h)^3}{6}$	$-a_{-1}\frac{(-h)^4}{24}$	$-a_{-1}\frac{(-h)^5}{120}$
$-a_0f_j$	$-a_0$	0	0	0	0	0
$-a_1f_{j+1}$	$-a_1$	$-a_1h$	$-a_1\frac{h^2}{2}$	$-a_1\frac{h^3}{6}$	$-a_1\frac{h^4}{24}$	$-a_1\frac{h^5}{120}$

- The algebraic system:

$$\begin{bmatrix} 0 & 0 & -1 & -1 & -1 \\ 1 & 1 & h & 0 & -h \\ -h & h & -h^2/2 & 0 & -h^2/2 \\ h^2/2 & h^2/2 & h^3/6 & 0 & -h^3/6 \\ -h^3/6 & h^3/6 & -h^4/24 & 0 & -h^4/24 \end{bmatrix} \begin{bmatrix} b_{-1} \\ b_1 \\ a_{-1} \\ a_0 \\ a_1 \end{bmatrix} = \begin{bmatrix} 0 \\ -1 \\ 0 \\ 0 \\ 0 \end{bmatrix} \implies \begin{bmatrix} b_{-1} \\ b_1 \\ a_{-1} \\ a_0 \\ a_1 \end{bmatrix} = \begin{bmatrix} 1/4 \\ 1/4 \\ 3/(4h) \\ 0 \\ -3/(4h) \end{bmatrix}$$

FINITE DIFFERENCE FORMULAE (XIII) — PADÉ APPROXIMATION

- The Padé approximation:

$$\frac{1}{4} \left(\frac{\delta f}{\delta x} \right)_{j+1} + \left(\frac{\delta f}{\delta x} \right)_j + \frac{1}{4} \left(\frac{\delta f}{\delta x} \right)_{j-1} = \frac{3}{4h} (f_{j+1} - f_{j-1})$$

Leading-order error $\frac{h^4}{30} f_j^{(v)}$ (**fourth-order accurate**)

- The approximation is **nonlocal**, in that it requires derivatives at the adjacent nodes which are also unknowns; Thus all derivatives must be determined at once via the solution of the following algebraic system

$$\begin{bmatrix} \ddots & \ddots & \ddots & & & \\ & 1/4 & 1 & 1/4 & & \\ & & \ddots & \ddots & \ddots & \\ & & & & & \ddots \end{bmatrix} \begin{bmatrix} \vdots \\ \left(\frac{\delta f}{\delta x} \right)_{j-1} \\ \left(\frac{\delta f}{\delta x} \right)_j \\ \left(\frac{\delta f}{\delta x} \right)_{j+1} \\ \vdots \end{bmatrix} = \begin{bmatrix} \vdots \\ \vdots \\ \frac{3}{4h} (f_{j+1} - f_{j-1}) \\ \vdots \\ \vdots \end{bmatrix}$$

- Closing the system at **endpoints** — use a lower-order one-sided (i.e., forward or backward) finite-difference formula

FINITE DIFFERENCE FORMULAE (XIV) — MODIFIED WAVENUMBER ANALYSIS

- Finite-Difference formulae applied to the Fourier mode $f(x) = e^{ikx}$ with the (exact) derivative $f'(x) = ik e^{ikx}$

- Central-Difference formula:

$$\left(\frac{\delta f}{\delta x}\right)_j = \frac{f_{j+1} - f_{j-1}}{2h} = \frac{e^{ik(x_j+h)} - e^{ik(x_j-h)}}{2h} = \frac{e^{ikh} - e^{-ikh}}{2h} e^{ikx_j} = i \frac{\sin(hk)}{h} f_j = ik' f_j,$$

where the **modified wavenumber** $k' \triangleq \frac{\sin(hk)}{h}$

- Comparison of the **modified wavenumber** k' with the **actual wavenumber** k shows how numerical differentiation errors affect different Fourier components of a given function

FINITE DIFFERENCE FORMULAE (XV) — MODIFIED WAVENUMBER ANALYSIS

- Fourth-order central difference formula

$$\begin{aligned} \left(\frac{\delta f}{\delta x}\right)_j &= \frac{-f_{j+2} + 8f_{j+1} - 8f_{j-1} + f_{j-2}}{12h} = \frac{2}{3h} (e^{ikh} - e^{-ikh}) f_j - \frac{1}{12h} (e^{ik2h} - e^{-ik2h}) f_j \\ &= i \left[\frac{4}{3h} \sin(hk) - \frac{1}{6h} \sin(2hk) \right] f_j = ik' f_j \end{aligned}$$

where the **modified wavenumber** $k' \triangleq \left[\frac{4}{3h} \sin(hk) - \frac{1}{6h} \sin(2hk) \right]$

- Fourth-order Padé scheme:

$$\frac{1}{4} \left(\frac{\delta f}{\delta x}\right)_{j+1} + \left(\frac{\delta f}{\delta x}\right)_j + \frac{1}{4} \left(\frac{\delta f}{\delta x}\right)_{j-1} = \frac{3}{4h} (f_{j+1} - f_{j-1}),$$

where $\left(\frac{\delta f}{\delta x}\right)_{j+1} = ik' e^{ikx_{j+1}} = ik' e^{ikh} f_j$ and $\left(\frac{\delta f}{\delta x}\right)_{j-1} = ik' e^{ikx_{j-1}} = ik' e^{-ikh} f_j$.

Thus

$$\begin{aligned} ik' \left(\frac{1}{4} e^{ikh} + 1 + \frac{1}{4} e^{-ikh} \right) f_j &= \frac{3}{4h} (e^{ikh} - e^{-ikh}) f_j \\ ik' \left(1 + \frac{1}{2} \cos(kh) \right) f_j &= i \frac{3}{2h} \sin(hk) f_j \implies k' \triangleq \frac{3 \sin(hk)}{2h + h \cos(hk)} \end{aligned}$$